

Spontaneous decay

Alessandro Cerè



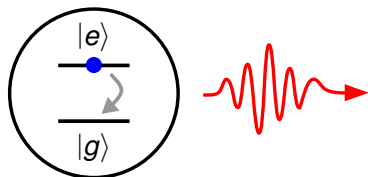
Outline

The problem

Wigner-Weisskopf Theory

A simple experiment

Spontaneous decay in free space



We consider a two level atom with ground and excited states respectively $|g\rangle$ and $|e\rangle$.

We will calculate the rate at which an atom spontaneously relaxes to a lower energy state while emitting a photon.

Initial state

$$|\psi(0)\rangle = |e, \{0\}\rangle$$

- the atom starts in excited state $|e\rangle$
- the radiation field consists of a continuum of modes which we think of as a cavity which is infinite in extent
- the radiation field starts in vacuum state $|\{0\}\rangle$

Final state

$$|\psi(\infty)\rangle = |g, 1_{\vec{k}}\rangle$$

The atom will eventually decay to the ground state $|g\rangle$ emitting a photon in mode \vec{k} .

Time dependent state

$$|\psi(t)\rangle = C_0^e(t)e^{-i\omega_0 t} |e, \{0\}\rangle + \sum_{\vec{k}} C_{1\vec{k}}^g(t)e^{-i\omega_k t} |g, 1_{\vec{k}}\rangle$$

Linear combination of the initial and final states provides a complete basis to describe the time-dependent state of the system.

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The Hamiltonian of the system consists of three parts

$$\hat{H} = \underbrace{\hbar\omega_0 |e\rangle\langle e|}_{H_A} + \underbrace{\sum_{\vec{k}} \hbar\omega_k \hat{a}_{\vec{k}}^\dagger \hat{a}_{\vec{k}}}_{H_F} + \underbrace{\sum_{\vec{k}} \hbar g_{\vec{k}} (|e\rangle\langle g| \hat{a}_{\vec{k}} + |g\rangle\langle e| \hat{a}_{\vec{k}}^\dagger)}_{H_I}$$

with atom-field coupling:

$$g_{\vec{k}} = i\sqrt{\frac{\omega_k}{2\hbar\epsilon_0 V}} (\vec{d} \cdot \vec{\epsilon}_{\vec{k}})$$

Time evolution of the system: Schrödinger equation

$$\hat{H}|\psi(t)\rangle = i\hbar \frac{\partial |\psi(t)\rangle}{\partial t}$$

Projecting onto $\langle e, \{0\} |$ and $\langle g, 1_{\vec{k}} |$

$$\frac{\partial}{\partial t} C_0^e(t) = i \sum_{\vec{k}} g_{\vec{k}} e^{-i(\omega_k - \omega_0)t} C_{1\vec{k}}^g(t) \quad (1)$$

$$\frac{\partial}{\partial t} C_{1\vec{k}}^g(t) = i g_{\vec{k}}^* e^{i(\omega_k - \omega_0)t} C_0^e(t) \quad (2)$$

Integration time!

We first formally integrate (2)

$$C_{1\vec{k}}^g(t) = ig_{\vec{k}}^* \int_0^t dt' e^{i(\omega_k - \omega_0)t'} C_0^e(t')$$

And then replace into (1)

$$\frac{\partial}{\partial t} C_0^e(t) = - \sum_{\vec{k}} |g_{\vec{k}}|^2 \int_0^t dt' e^{i(\omega_k - \omega_0)(t-t')} C_0^e(t')$$

From discrete to continuum

For $V \rightarrow \infty$ the discrete set of modes tends to a continuum

$$\sum_{\vec{k}} \rightarrow \overbrace{2}^{\text{2 pol per mode}} \int d^3k \underbrace{D(k)}_{\frac{V}{(2\pi)^3}}$$

$$\sum_{\vec{k}} \rightarrow \frac{2V}{(2\pi)^3} \int_0^{2\pi} d\phi \int_0^\pi \sin \theta d\theta \int_0^\infty k^2 dk$$

Let's get rid of $\sum_{\vec{k}} |g_{\vec{k}}|^2$

$$\sum_{\vec{k}} |g_{\vec{k}}|^2 \rightarrow \int_0^\infty dk k^2 \frac{\omega_k}{(2\pi)^3 \hbar \epsilon_0} \underbrace{\left[\int_0^\pi \sin \theta d\theta \int_0^{2\pi} d\phi \left(\vec{d} \cdot \vec{\epsilon}_{\vec{k}} \right)^2 \right]}_{\frac{8\pi}{3} |\vec{d}|^2}$$

using $\omega_k = ck$:

$$\sum_{\vec{k}} |g_{\vec{k}}|^2 \rightarrow \frac{|\vec{d}|^2}{6\pi^2 \hbar \epsilon_0 c^3} \int_0^\infty \omega_k^3 d\omega_k$$

Let's go back to the time integral

$$\int_0^t dt' e^{i(\omega_k - \omega_0)(t-t')} C_0^e(t')$$

Wigner-Weisskopf approximation

We assume that $C_0^e(t)$ varies in a timescale $\Gamma \ll \omega_0$

$$C_0^e(t') \rightarrow C_0^e(t)$$

And we can take it out of the integral

$$C_0^e(t) \int_0^t dt' e^{i(\omega_k - \omega_0)(t-t')}$$

This is a Markov approximation:
the system has no memory of the past.

Time integration

We're interested in time scales $t \gg 1/\omega_0$, then $\int_0^t \rightarrow \int_0^\infty$

$$\int_0^\infty dt' e^{i(\omega_k - \omega_0)(t-t')} = \pi\delta(\omega_k - \omega) - iP\left(\frac{1}{\omega_k - \omega}\right)$$

where P is Cauchy principal part.

Complex values correspond to frequency shifts (Lamb shift).
It's nasty, it diverges. Let's forget about it...

Combining temporal and momentum part

$$\frac{\partial}{\partial t} C_0^e(t) = - \underbrace{\sum_{\vec{k}} |g_{\vec{k}}|^2}_{\text{momentum part}} \underbrace{\int_0^t dt' e^{i(\omega_k - \omega_0)(t-t')} C_0^e(t')}_{\text{temporal part}}$$

$$\frac{\partial}{\partial t} C_0^e(t) = \overbrace{\frac{|\vec{d}|^2}{6\pi^2 \hbar \epsilon_0 c^3}}^{\text{momentum part}} \overbrace{\int_0^\infty d\omega_k \omega_k^3 \pi \delta(\omega_k - \omega_0) C_0^e(t)}^{\text{temporal part}}$$

Excited state amplitude time evolution

$$\frac{\partial C_0^e(t)}{\partial t} = -\frac{\Gamma}{2} C_0^e(0) \quad \text{with} \quad \Gamma = \frac{|\vec{d}|^2 \omega_0^3}{3\pi\hbar\epsilon_0 c^3}$$

this one with can solve:

$$C_0^e(t) = e^{-\frac{\Gamma}{2}t} C_0^e(0) \tag{3}$$

The excited state amplitude decays exponentially with rate Γ ,
AKA Einstein A coefficient

Spectrum of the emitted light

We can now solve Eq. (2)

$$\frac{\partial}{\partial t} C_{1\vec{k}}^g(t) = ig_{\vec{k}}^* e^{i(\omega_k - \omega_0)t} e^{-\frac{\Gamma}{2}t} C_0^e(0)$$

The process takes an infinite time to complete. Integrating:

$$\lim_{t \rightarrow \infty} \left| C_{1\vec{k}}^g(t) \right| = \frac{|ig_{\vec{k}}|}{\frac{\Gamma^2}{4} + (\omega_k - \omega_0)^2}$$

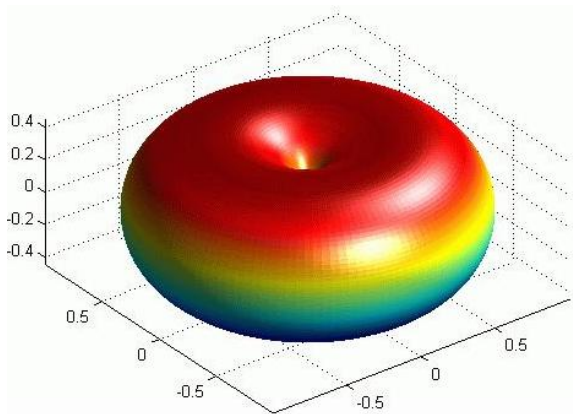
We find a Lorentzian spectrum.

What a surprise, eh?

Spatial distribution

The $|ig_{\vec{k}}|$ term contains the spatial information. For $\Delta m = 0$

$$\lim_{t \rightarrow \infty} |C_{1\vec{k}}^g(t)| \sim \sum_s |\vec{d} \cdot \vec{\epsilon}|^2 \propto \sin^2(\theta)$$



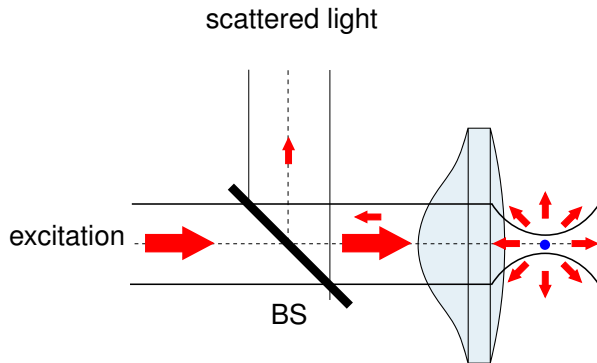
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The problem

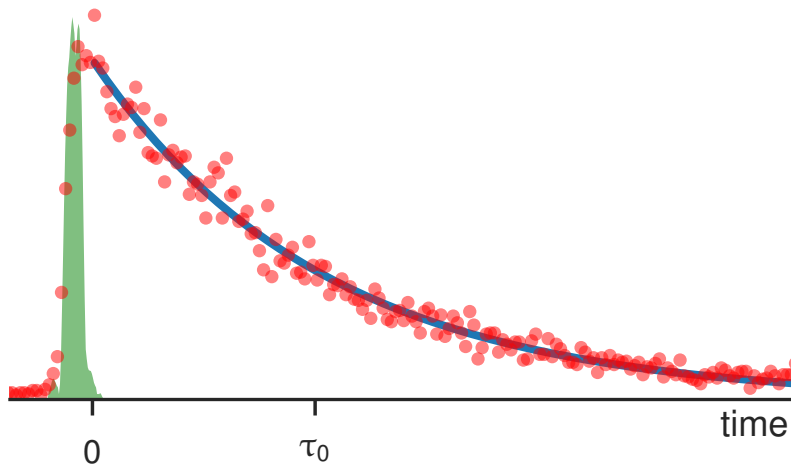
Wigner-Weisskopf Theory

A simple experiment

Pulse excitation of a single atom



Spontaneous decay of ^{87}Rb



cqtac@nus.edu.sg
qolah.org
quantumlah.org

